

An inhomogeneous and anisotropic constitutive model of human dentin

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Abstract

Dentin constitutes the major part of human tooth. It is composed of a large number of tubules with both variational radii and radially parallel pattern. In addition, peritubular dentin surrounds each tubule lumen and has a higher elastic modulus than the matrix of dentin, i.e. intertubular dentin. Considering the above microstructural characteristics, a micromechanics model is used in this paper to evaluate the overall elastic properties of dentin. Five independent effective elastic parameters in transverse isotropic elasticity matrix can be expressed analytically by the material parameters of peri- and intertubular dentin and the volume fraction of tubules. To determine the effectivity of this theoretical model, a finite element (FE) model simulating a longitudinal tooth slice in moiré fringe testing of Wang and Weiner (*J. Biomech.* 31 (1998) 135) was performed. Furthermore, the FE model was developed incorporate modeling of variation of tubule's diameter and softer characteristic of intertubular dentin near the dentin–enamel junction and around the pulp chamber. It turned out that the isoline figure of longitudinal displacement by FE calculation has very similar patterns to the moiré fringe results. However, the FE results of displacement by traditional stress–strain models which regard dentin as a homogeneous and isotropic material show an obviously different strain distributions as compared to published moiré fringes results. Thus the inhomogeneous and anisotropic model developed in this paper more accurately reflects the true physical nature of human dentin.

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1. Introduction

In dentin, fluid-filled microscopic tubules extend through the entire dentin from the pulp wall to the dentin–enamel junction (DEJ) or the cementum. In the longitudinal section of dentin (*Fig. 1a*), the tubules in the root of tooth are approximately parallel, but the tubules in the crown of tooth are radial. In the transverse section (*Fig. 1b*), the tubules are also radially oriented. The peri- and intertubular dentin represent the mineralized portions in dentin. The peritubular dentin encircles the tubules and is characterized by its high mineral content (*Mjör and Fejerskov, 1986*). The intertubular dentin occupies the volume outside of peritubular dentin and contains more collagen and less mineral than the peritubular dentin (*Fig. 1c*).

As dentin is examined closer to the pulp cavity, the diameters of tubules increase gradually, and moreover, the numerary density of tubules also has an increasing trend, although the total number of tubules remain constant (*Ketterl, 1961; Tronstad, 1973; Garberoglio and Brännström, 1976; Schilke et al., 2000*). The above structural characters of dentin should lead to an inhomogeneous and anisotropic stress–strain relation. In fact, the effects of tubules on the mechanical properties of dentin were considered in the pioneer work of *Peyton et al. (1952)*, but no effects were observed due to the unavailability of micromechanical methods at that time. This led to dentin being treated as a homogeneous and isotropic elastic material.

More recently, several groups have recognized that the presence of dentinal tubules may influence the mechanical properties of dental tissues, e.g. bond strength of restoration (*Phrukkanon et al., 1999; Ogata et al., 2001*). Microhardness (*Pashley et al., 1985; Meredith et al., 1996; Wang and Weiner, 1998*) and

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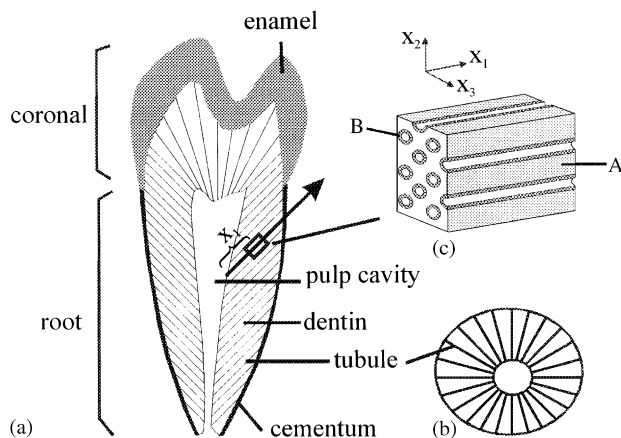


Fig. 1. Schematic figures of a premolar and the microstructure of its dentin: (a) the longitudinal section of the premolar; (b) the transverse section of the dentin; (c) a representative volume element (RVE). A—intertubular dentin, B—peritubular dentin.

nanoindentation (Kinney et al., 1996, 1999) have been adopted to determine the microstructure and physical properties of dentin. Others have reported changes in different chemical components of dentin along the tubules (Wentrup-Byrne et al., 1997).

Furthermore, Kinney et al. (1999) and Huo and Zheng (1999) created micromechanics models of dentin. In Kinney's model, the effects of peri- and intertubular dentin on anisotropic properties of dentin were considered. They measured the elastic modulus of the peritubular dentin as 30 GPa and that of intertubular dentin as 15 GPa. They presented a theoretical model in which the tubule lumen lined by a peritubular dentin cuff was assumed to be an isotropic cylindrical inclusion. That assumption may not be appropriate if the composite made up of tubule and peritubular dentin behaves as anisotropic material.

In the authors' previous work (Huo and Zheng, 1999), a varying transverse isotropic stress–strain relation of dentin was established. That theoretical model can reflect the effect of tubules on the inhomogeneous and anisotropic properties of dentin. In that model, however, the peri- and intertubular dentine were regarded as materials with same elastic properties. In this paper the previous model is extended by considering the different mechanical properties of peri- and intertubular dentin, and then a numerical simulation is presented and compared with the previously published moiré fringe testing results (Wang and Weiner, 1998).

2. Materials and methods

2.1. Theoretical model

A representative volume element (RVE) was taken out from the dentin (see Fig. 1c). The RVE was small

enough so that the dentin tubules in it could be regarded as being parallel and diameters of the tubules as being invariant. The RVE should also be large enough to contain many tubules so that it can represent the transversely isotropic dentin material in both the local and macroscopic sense. A Cartesian coordinate system $\{x_1, x_2, x_3\}$ was used, with the x_1 -axis along the dentin tubule direction, the privileged direction of the transverse isotropy.

Using the one-index notations $\tau_1 = \sigma_{11}$, $\tau_4 = \sigma_{23}$, $\gamma_2 = \varepsilon_{22}$, $\gamma_5 = 2\varepsilon_{31}$, etc., instead of the usual second-order symmetric stress tensor σ_{ij} and strain tensor ε_{ij} , the overall elastic stress–strain relationship can be written into the matrix form

$$\gamma = \mathbf{D}\tau, \quad (1)$$

where the components of elasticity compliance matrix \mathbf{D} can be expressed as

$$\begin{aligned} D_{11} &= 1/E', & D_{22} &= D_{33} = 1/E, \\ D_{44} &= (1 + \nu)/E, & D_{55} &= D_{66} = 1/2G', \\ D_{12} &= D_{21} = D_{13} = D_{31} = -\nu'/E', \\ D_{23} &= D_{32} = -\nu/E, \end{aligned} \quad (2)$$

other components vanish. The E , ν , E' , ν' and G' are the five independent effective elastic parameters of the transversely isotropic material. E' is the Young's modulus along the direction of tubule. E and ν are the Young's modulus and Poisson's ratio within the isotropic plane, respectively. G' and ν' are the shear modulus and Poisson's ratio in the plane of x_1x_2 or x_1x_3 , respectively.

Assuming that the peri- and intertubular dentin are homogeneous and isotropic, corresponding elasticities are denoted by \mathbf{C}^p and \mathbf{C}^i , Young's moduli by E_p and E_i , Poisson's ratios by ν_p and ν_i , respectively. The transverse shapes of tubule and peritubular dentin are supposed to be similar and coaxial circles and their transverse areas within the RVE are denoted by A_t and A_p , respectively. The transverse area of the intertubular dentin within the RVE is denoted by A_i .

When an ellipsoidal inclusion which contains another ellipsoidal inclusion is embedded in an infinite matrix material, the double-inclusion method can present the overall elasticity of the double inclusions (Nemat-Nasser and Hori, 1993, p. 352). Here a tubule and its surrounding peritubular dentin are considered as the double inclusions, which are assumed to be ellipsoidal with an infinitely long axis parallel to the tubules in RVE. It is considered that in RVE (Fig. 1c) an inclusion (peritubular dentin) including another cavity (dentin tubule) is embedded in matrix material (intertubular dentin). So the overall elasticity of the tubule–peritubule composite is given by

$$\mathbf{C}^{\text{tp}} = \mathbf{C}^i[\mathbf{I} + (\mathbf{S} - \mathbf{I})\mathbf{A}](\mathbf{I} + \mathbf{S}\mathbf{A})^{-1}, \quad (3)$$

where the matrix form of Eshelby’s tensor \mathbf{S} is given with their components,

$$\begin{aligned} S_{21} = S_{31} &= 4v_i/P, & S_{23} = S_{32} &= (4v_i - 1)/P, \\ S_{22} = S_{33} &= (5 - 4v_i)/P, & S_{44} &= 2(3 - 4v_i)/P, \\ S_{55} = S_{66} &= 4(1 - v_i)/P, \end{aligned} \quad (4)$$

other components vanish, in which $P = 8(1 - v_i)$. Eshelby’s tensor is an important fourth-order tensor in the averaging scheme of micromechanics and it reflects the effect of the eigenstrain within an inclusion with ellipsoidal shape on its surrounding matrix. This tensor depends on the shape of inclusion and the properties of the matrix materials and is independent of the material properties of inclusion. Here the tubule and its surrounding peritubular dentin are assumed to be concentric circles, so their Eshelby’s tensors are same.

In (3), \mathbf{A} is defined by

$$\mathbf{A} = f_{tp}(\mathbf{I} - \mathbf{S})^{-1} + (1 - f_{tp})[(\mathbf{C}^i - \mathbf{C}^p)^{-1}\mathbf{C}^i - \mathbf{S}]^{-1}. \quad (5)$$

In the above expression, $f_{tp} = A_t/(A_t + A_p)$ and \mathbf{I} denotes the unit matrix. Note that the elasticity of tubules is null in definition (5).

When the elasticity of the matrix deviates considerably from that of inclusion, the two-phase model can be used to evaluate the interaction between inclusions and matrix material (Nemat-Nasser and Hori, 1993, p. 343). In this model, a typical inclusion is embedded in a finite ellipsoidal region with matrix elasticity and then this double-inclusion is embedded in an infinite solid. The volume fraction of the inclusion in that subregion is equal to the actual one of the inclusions in total materials. So the average response of the above subregion can be given to the overall moduli of the two-phase material. Now consider the tubule–peritubule composite as one inclusion with elasticity \mathbf{C}^{tp} which is embedded in the matrix material of intertubular dentin, then the overall elasticity of dentin can be given as follows:

$$\mathbf{C}^d = \mathbf{C}^i \{ \mathbf{I} + f_p(\mathbf{S} - \mathbf{I})[(\mathbf{C}^i - \mathbf{C}^{tp})^{-1}\mathbf{C}^i - \mathbf{S}]^{-1} \} \{ \mathbf{I} + f_p\mathbf{S}[(\mathbf{C}^i - \mathbf{C}^{tp})^{-1}\mathbf{C}^i - \mathbf{S}]^{-1} \}^{-1}, \quad (6)$$

where $f_p = (A_t + A_p)/(A_t + A_p + A_i)$. So the volume fraction of tubule is $f_t = A_t/(A_t + A_p + A_i) = f_{tp} * f_p$.

Substituting (2)–(5) into (6) and solving the following matrix equation:

$$\mathbf{D} = (\mathbf{C}^d)^{-1}, \quad (7)$$

the five independent elastic parameters, E , ν , E' , ν' and G' , can be expressed by the material parameters of E_p , E_i , ν_p , ν_i , f_{tp} , f_p . Note that during the above derivations, the inclusions within RVE are implied to be cylindrical, but in fact the volume fractions f_{tp} and f_p vary with the coordinate x_1 , i.e. $f_{tp} = f_{tp}(x_1)$, $f_p = f_p(x_1)$. In this paper, the above expressions for RVE are assumed to still be valid for whole dentin. The overall compliance matrix \mathbf{D}

of dentin is dependent on x_1 and the stress–strain relationship is *varying* transverse isotropic.

If the peri- and intertubular dentin are regarded as the same material, i.e. $E_p = E_i = E_m$, $\nu_p = \nu_i = \nu_m$, $f_{tp}(x_1) = 1$, $f_p(x_1) = f_t(x_1) = c(x_1)$, the expressions of E , ν , E' , ν' and G' coincide with the previous results by Huo and Zheng (1999).

2.2. The establishment of a 2D finite element model

Wang and Weiner (1998) cut a bucco-lingual section 2 mm thick out of the central part of a human canine, and applied a compressive force of 500 N to the top of section. By means of a moiré fringe technique, they presented the displacement distribution, which will be referred to as *Wang–Weiner image* in this paper, and calculated the strain distribution in the section. As their experiment was simple and was easily simulated, a finite element model was established with the same geometric configuration as their work (Fig. 2a). The DEJ in this FE model had same path as in the model of Wang and Weiner (Fig. 5b). The bottom of the FE model was fixed. A load of magnitude 500 N was uniformly applied at the top of the model. The FE mesh was generated with 4064 elements and 4183 nodes. All of the preprocessing, calculating and postprocessing were done with the software MSC.Patran 2003.

As the diameter of a human tooth is about 8 mm, a 2 mm thick slice may be too thick for assuming a 2D structure. However, building a 3D FEM model with inhomogeneous and anisotropic constitutive model requires adequate geometric data in the thickness direction of slice, e.g. the orientation of tubules and the position of DEJ, which entail more assumptions. Considering that the tubules on the front and back

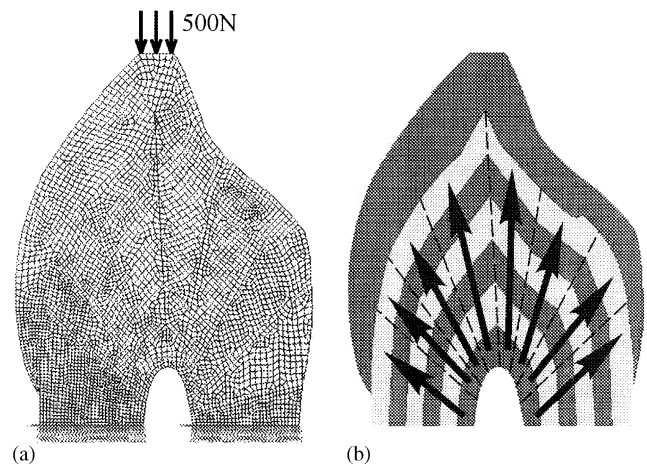


Fig. 2. (a) The finite element grid and (b) the definitions of inhomogeneous and anisotropic material properties of the section of a human tooth. The dentin is uniformly partitioned into six layers along the tubules. The arrows represent the dentin tubule direction, i.e. the privileged direction of the transverse isotropy.

planes of slice are symmetrical to the mid-plane and their angle diversity is not very great, it will not affect the computation results greatly to neglect the geometric information in the slice's thickness.

Four FE models with different properties of dentin were established for comparison with *Wang–Weiner image*. The material properties of enamel for all FE models were assigned a Young's modulus of 80 GPa and a Poisson's ratio of 0.3 (Rees and Jacobsen, 1997). Most researchers ignore the effect of tubules on the mechanical properties of dentin and regard dentin as homogeneous and isotropic, that model will be called the *HI Model* in this paper. In the *HI Model*, the elastic modulus was assigned as 15 GPa and Poisson's ratio as 0.31 (Rees and Jacobsen, 1997).

In the following, it will be described how to establish an inhomogeneous and anisotropic FE model of dentin.

To reflect the variation in mechanical properties, the elements of dentin were equally partitioned into six layers along the tubules and these layers were assigned different material properties according to the above inhomogeneous and anisotropic (IA) stress–strain relation (Fig. 2b).

Fig. 3 shows the curves of three volume fractions, $f_{tp}(x_1)$, $f_t(x_1)$ and $f_p(x_1)$, varying with the normalized distance, x_1/L , where L is the length of tubule. Garberoglio and Brännström (1976) measured the diameter and numerical density of tubules in different locations in dentin. Based on their results, the values of $f_t(x_1)$ were calculated. Note that at the wall of the pulp chamber, the value, $f_t(0) = 0.32$, was directly measured from the SEM photograph by Garberoglio and Brännström (1976). Pashley (1989) assumed the area occupied by peritubular dentin as three times that of the tubules, so $f_p(x_1) = 4f_t(x_1)$. Within the predentin, which is close to the pulp cavity, however, there is only intertubular dentin and no peritubular dentin (Mjör and

Fejerskov, 1986), i.e. $f_p(0) = f_t(0)$. Therefore, the volume fraction $f_p(x_1)$ may be overestimated when x_1 approaches zero. Now the area ratio of peritubular dentin was directly measured, i.e. $f_p(0.5 \text{ mm}) = 0.31$, based on the SEM photograph after decalcification at the position of 0.5 mm from the pulp wall (Garberoglio and Brännström, 1976). The values of $f_{tp}(x_1)$ can be calculated by the equation, $f_t(x_1) = f_{tp}(x_1) * f_p(x_1)$.

Kinney et al. (1999) measured the Young's moduli of peri- and intertubular dentin as 30 and 15 GPa, respectively. From the nanoindentation testing, Kinney et al. (1996) showed that the intertubular dentin became soft approaching either the pulp cavity or the DEJ, but the hardness of the peritubular dentin remained constant. This led us to assume the Young's modulus of intertubular dentin in layers 1 and 2 towards the pulp cavity as 5 and 10 GPa and in layer 6 near the DEJ as 10 GPa, respectively. As there were no experimental values for Poisson's ratio of peri- and intertubular dentin, they were assumed to be 0.3 and 0.4, respectively. Poisson's ratio of intertubular dentin was assigned the larger value because there are more organic components in it, e.g. collagen fibrils.

The six material parameters of every layer of dentin, E_p , E_i , v_p , v_i , f_{tp} , f_p , could be obtained through the above procedure and then substituted into Eq. (7) to compute the corresponding effective material parameters. This FE model is called *IA Model* in this paper.

Assuming that the peri- and intertubular dentine were the same matrix material, Huo and Zheng (1999) presented a varying transverse isotropic stress–strain relation of dentin. Based on this theoretical model, the elastic modulus of matrix material was concluded as 29.5 GPa and Poisson's ratio as 0.44 from their experiments (Huo et al., 2000). The effective material parameters of dentin can be calculated by the theoretical model of Huo and Zheng (1999). This FE model is called *Huo–Zheng Model* in this paper.

Kinney et al. (1999) presented a theoretical model in which the tubule lumen lined by a peritubular dentin cuff was assumed to be an isotropic cylindrical inclusion. In their paper, the elastic modulus of inclusion, E_f , was calculated as 22.5 GPa and Poisson's ratio, v_f , was assumed to be 0.25. The elastic modulus of matrix (i.e. intertubular dentin), E_m , was measured as 15 GPa and Poisson's ratio, v_m , was assumed to be 0.4. Substituting the above parameters into their theoretical model, the effective material parameters of dentin can be calculated. This fourth FE model is called the *Kinney Model* in this paper.

In the above four FE models, the *IA Model* considered all the features of dentin's mechanical properties based on the presently published experimental data. The *Huo–Zheng Model* and the *Kinney Model* were both previously published where the former neglected the difference of peri- and intertubular dentin

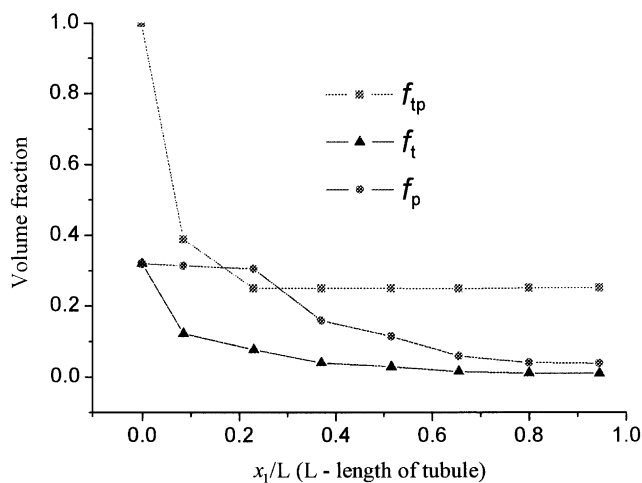


Fig. 3. The curves of three volume fractions f_{tp} , f_t and f_p , varying with the normalized distance x_1/L (refer to Garberoglio and Brännström, 1976; Pashley, 1989).

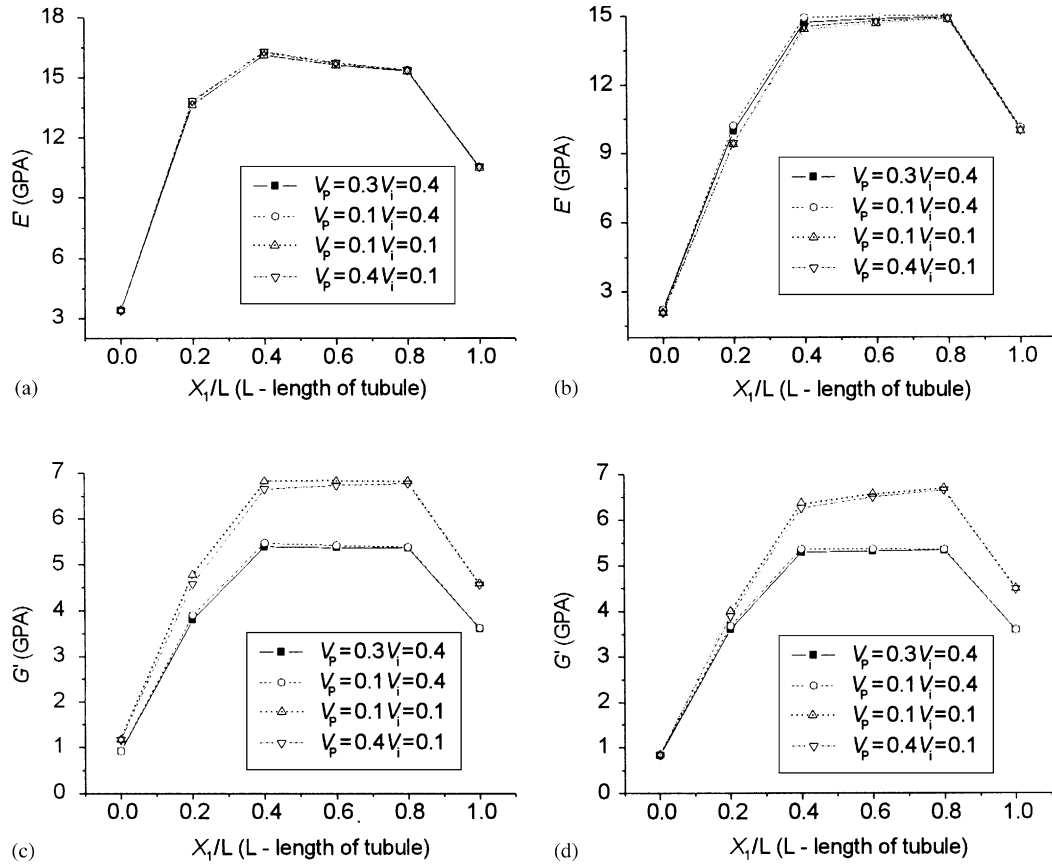


Fig. 4. The effective elastic parameters of dentin at the different normalized distance x_1/L for the varying Poisson's ratios of peri- and intertubular dentin, v_p and v_i . E' —Young's modulus in the direction of tubules; G' —shear modulus in the plane of x_1x_2 or x_1x_3 ; E —Young's modulus within the isotropic plane; G —shear modulus within the isotropic plane ($= E/[2(1 + \nu)]$).

and the latter simply assumed the lumen surrounded by peritubular dentin as an isotropic inclusion. Now people usually use the so-called *HI Model*, which does not consider the microstructure of dentin at all.

3. Results

Poisson's ratios of peri- and intertubular dentin, v_p and v_i , cannot be measured directly, so its effect on the effective elastic parameters of dentin should be examined. Four pairs of v_p and v_i were presented, i.e. 0.3 and 0.4, 0.1 and 0.4, 0.1 and 0.1, 0.4 and 0.1. The first pair had been assigned to the *IA FE Model*. Other pairs represented some combinations of v_p and v_i with larger difference. All of them were substituted into Eq. (7) to calculate the effective elastic parameters of dentin at the different distances from the pulp wall. It can be found that all the elastic parameters have a similar trend, i.e. being lowest near the pulp wall and then increasing up to a platform, finally decreasing near the DEJ (Fig. 4). This variation trend is very similar to the experimental hardness result of Wang and Weiner (1998). It is shown that Young's moduli E and E' are negligibly affected by the different cases of Poisson's ratios (Figs. 4a and b).

Little effect of Poisson's ratio v_p of peritubular dentin on shear moduli G' and G ($= E/[2(1 + \nu)]$) was found (Figs. 4c and d). However, the shear moduli obviously change for varying Poisson's ratio v_i of intertubular dentin at some positions with a maximum relative difference of them about 25%. The effective Poisson's ratio ν' in the plane of x_1x_2 or x_1x_3 is also dominated by v_i and it barely varies at the distance from the pulp wall, which is not shown in the figure.

Fig. 5a is a copy of the fringe pattern of the longitudinal displacement field from the *Wang–Weiner image*. According to the two FE models, *HI Model* and *IA Model*, the isoline images of longitudinal displacement are shown in Figs. 5b and c. The corresponding results of two micromechanics model, *Huo–Zheng Model* and *Kinney Model*, are presented in Figs. 6a and b. To compare the distribution pattern of longitudinal displacement computed by the different models of dentin with the fringe image, the contour lines within the dentin were plotted together (see Figs. 7a and b).

When using the *HI Model*, it was predicted that the inner part of dentin had the same ability of transferring incisal load as the outer part (Fig. 5b). Thus, the distribution pattern reveals a very flat form in most regions of dentin (Fig. 7a). Only when approaching the



Fig. 5. (a) The fringe figure of the *Wang–Weiner image* (Wang and Weiner, 1998). The calculated compressive strain values (%), the negative sign for the compressive strain was omitted) at different locations of the sample was given by Wang and Weiner (1998). The tip of tooth crown, which was not displayed in the original paper, is redrawn by the dash line. The isoline images of longitudinal displacement are calculated by (b) *HI Model* and (c) *IA Model* of dentin.

pulp cavity, is there a slight curving shape for the isoline due to the effect of the pulp chamber.

The *IA Model* considered the more porous character of dentin near the pulp wall and was constructed so that the intertubular dentin near the DEJ and the pulp wall had softer material properties. Approaching pulp cavity, the *IA Model* predicted obvious curving patterns for longitudinal displacement isolines within dentin (Fig. 5c). Furthermore, the displacement isolines for the *IA Model* had a slight turn downwards in the dentin near the DEJ. It can be found that the FE analysis for the *IA*

Model showed very similar fringe patterns to *Wang–Weiner image* (Fig. 7a). These characters of isoline pattern were not displayed by the *HI Model*.

When the stress–strain relation of dentin by Huo and Zheng (1999) was used (Fig. 6a), the displacement field had similar curving patterns to *Wang–Weiner image* within the center part of dentin (Fig. 7b); but a turn downwards of displacement isolines near the DEJ was not predicted since the softer properties of intertubular dentin near the DEJ was not considered. Being similar to *Huo–Zheng Model*, the *Kinney Model* could not

predict the turn downwards near the DEJ (7b). It was significantly distinguished from *Wang–Weiner image* in that the longitudinal displacement field of the *Kinney Model* had a flat form in most regions of dentin.

The above FE results show that the *IA Model* most closely matched the physical properties of dentin shown by *Wang–Weiner image*. Other models, such as *HI*, *Huo–Zheng* and *Kinney* ones, ignored some features of the mechanical properties of dentin more and less. Therefore, these models gave different patterns of longitudinal displacement from the moiré fringe image.

4. Discussions

Two averaging schemes, i.e. double-inclusion method and two-phase model, are used in this paper. An analytical expression of the overall elasticity of dentin is obtained, but the solved expressions of the elastic parameters are too long to be written out integrally in this paper. In fact, the former method is the generalized one of the latter. In the two schemes, the mathematical derivation is exact. Interaction between the inclusions and their immediate surrounding matrix material of

different elasticity is directly included in the models (Nemat-Nasser and Hori, 1993, p. 340). The outmost region, however, is assumed to be an infinite matrix material for these two methods, so they may magnify the overall elasticity of the composite when the volume fraction of tubules is large. Therefore, the curvature of displacement isolines close to the pulp cavity in the actual case may be greater than that in the computation results of this paper.

By this inhomogeneous and anisotropic theoretical model of dentin, the FE analysis results of longitudinal displacement were very similar to a previous moiré fringe results in two aspects, i.e. highly curving close to the pulp cavity within the central section of dentin and slightly turning downwards near the DEJ. There are two characters of material properties near the pulp cavity, i.e. being more porous and softer, which contribute to the curving patterns there. Ketterl (1961) measured the value of 76% for incisors, while Schug-Kösters and Ketterl (1973) obtained a larger value of 79%. However, several reports gave much smaller volume fractions of 14% to 32% (Tronstad, 1973; Garberoglio and Brönström, 1976; Schilke et al., 2000). The width of collagen and proteoglycan-rich predentin matrix at the boundary of the pulp chamber is thought to be about 50 μm thick. So it is believed that even if there are higher volume fractions of tubules near the pulp chamber, their influence fall off rapidly away from the pulp wall, due to the need of enduring the occlusal loading. When establishing the *IA FE Model*, the elastic modulus of intertubular dentin was defined as smaller values, up to 5 GPa close to the pulp wall and 10 GPa near the DEJ. The softer material properties of intertubular dentin in the above regions contribute to the curving patterns near the pulp cavity and lead to turning downwards of displacement isolines near the DEJ. When the softer properties of peri- and intertubular dentin were ignored, i.e. E_p and E_i were assumed to be constant, a turn downwards of displacement isolines near the DEJ disappeared and the curvatures of them towards the

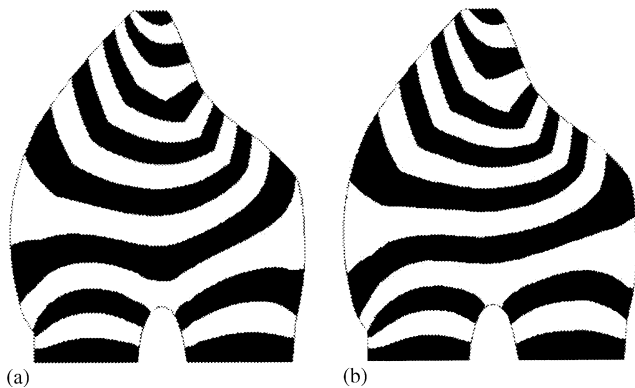


Fig. 6. The isoline images of longitudinal displacement are calculated by (a) *Huo–Zheng Model* and (b) *Kinney Model* of dentin.

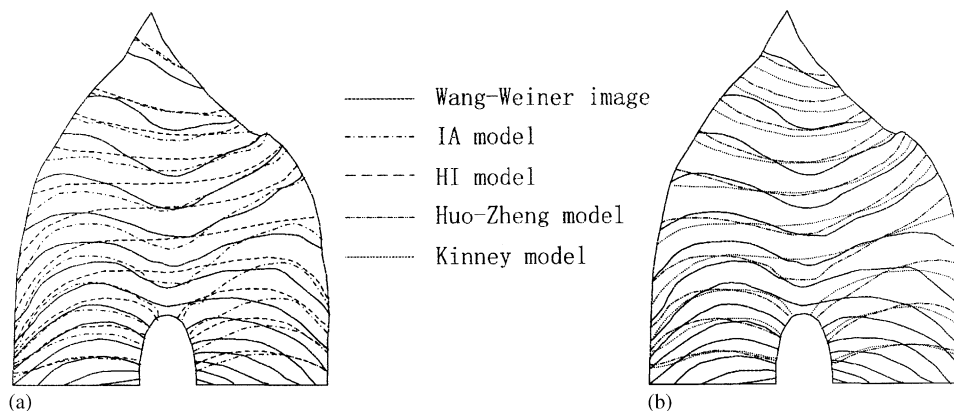


Fig. 7. The comparison of the contour lines of longitudinal displacement within the dentin computed by different FE models with the ones obtained from the *Wang–Weiner image*.

pulp cavity became less. So the inhomogeneity has greater impact to the distribution pattern than the anisotropy (also refer to Huo and Zheng, 2002).

The experimental data for the Young's moduli of peri- and intertubular dentin along the thickness of dentin are still imperfect. Some macroscopic tests had been done to study the hardness of dentin from the DEJ towards the pulp cavity (Smith and Cooper, 1971; Pashley et al., 1985; Watanabe et al., 1996; Lertchirakarn et al., 2001). Using these corresponding methods, the mechanical properties of peri- and intertubular dentin could not be separately determined. Owing to the development of atomic force microscopy fitted with a diamond stylus, however, one can look at the mechanical properties of intertubular dentin separated from peritubular dentin. Some studies showed that the Young's modulus and hardness of peritubular dentin did not vary with depth and the hardness of intertubular dentin in DEJ was higher than near the pulp (Kinney et al., 1999; Marshall et al., 2001). But more detailed testing for the Young's moduli along the tubule's direction may still be necessary.

Poisson's ratios of peri- and intertubular dentin have not yet been experimentally determined because of their small size, but an experiment may be designed by the theoretical model developed in this paper. For example, a small cube extracted from dentin can be compressed parallel and perpendicular to the tubule directions by a specially modified atomic-force microscope and the corresponding load–displacement measurements can be easily obtained. This would permit two Poisson's ratios to be obtained by solving the two parameter equations along the directions parallel and perpendicular to the tubule. The use of accurate Poisson's ratio would improve the FE simulations.

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